

3D → 2D GPE Reduction

let us start with fully dimensional 3D GPE eq:

$$i\hbar \psi_t = \left[-\frac{\hbar^2}{2m} \nabla^2 + \frac{m}{2} \underbrace{(\omega_x^2 x^2 + \omega_y^2 y^2 + \omega_z^2 z^2)}_{V(x,y,z)} + g |\psi|^2 \right] \psi \quad (1)$$

We suppose a "pancake" trap where $\omega_z \gg \omega_x \approx \omega_y \equiv \omega_r$
 $\Rightarrow V(x,y,z) \rightarrow \frac{m}{2} (\omega_r^2 r^2 + \omega_z^2 z^2)$

In this quasi-2D scenario the wave function will not be able to be excited in the z -direction $\Rightarrow$ in the z -direction the system will be in its "ground state".

[ground state of a harmonic trap $\rightarrow$ quantum harmonic oscillator
 $\rightarrow$ Gaussian]

- ∴ let us do :
- a) separation of variables : $\psi_{3D} = \psi_{2D} \cdot \phi_0(z) \cdot f(t)$
 - b) method of averaging ~~is~~ over z -direction

3D → 2D : $\psi_{3D} = \psi_{2D}(x,y) \phi_0(z) e^{-i\frac{\delta}{\hbar} t}$ (separation of variables)

$$\Rightarrow \phi_0 \left[i\hbar \frac{\partial}{\partial t} \psi_{2D} e^{-i\frac{\delta}{\hbar} t} \right] = \int_z \left(-\frac{\hbar^2}{2m} \nabla_{2D}^2 \psi_{2D} \right) \phi_0 e^{-i\frac{\delta}{\hbar} t} + \int_z \left[\frac{\hbar^2}{2m} \phi_0'' e^{-i\frac{\delta}{\hbar} t} + [V + g \phi_0^2 |\psi_{2D}|^2] \phi_0 \psi_{2D} e^{-i\frac{\delta}{\hbar} t} \right]$$

$$\Rightarrow i\hbar \frac{\partial}{\partial t} [\psi_{2D} e^{-i\frac{\delta}{\hbar} t}] = -\frac{\hbar^2}{2m} \nabla_{2D}^2 \psi_{2D} e^{-i\frac{\delta}{\hbar} t} - \frac{\hbar^2}{2m} \frac{\phi_0''}{\phi_0} \psi_{2D} e^{-i\frac{\delta}{\hbar} t} + [V + g \phi_0^2 |\psi_{2D}|^2] \psi_{2D} e^{-i\frac{\delta}{\hbar} t}$$

$$\Rightarrow i\hbar \left(\frac{\partial}{\partial t} \psi_{2D} \right) e^{-i\frac{\delta}{\hbar} t} + i\hbar \psi_{2D} \left(-i\frac{\delta}{\hbar} \right) e^{-i\frac{\delta}{\hbar} t} = -\frac{\hbar^2}{2m} \nabla_{2D}^2 \psi_{2D} e^{-i\frac{\delta}{\hbar} t} - \frac{\hbar^2}{2m} \frac{\phi_0''}{\phi_0} \psi_{2D} e^{-i\frac{\delta}{\hbar} t} + [V + g \phi_0^2 |\psi_{2D}|^2] \psi_{2D} e^{-i\frac{\delta}{\hbar} t}$$

$$\div \psi_{2D} \Rightarrow i\hbar \frac{\partial \psi_{2D}}{\partial t} + \frac{\hbar^2}{2m} \frac{\nabla_{2D}^2 \psi_{2D}}{\psi_{2D}} = -\frac{\hbar^2}{2m} \frac{\phi_0''}{\phi_0} + [V + g \phi_0^2 |\psi_{2D}|^2]$$

$$i\hbar \frac{\partial \psi_{2D}}{\partial t} + \frac{\hbar^2}{2m} \frac{\nabla_{2D}^2 \psi_{2D}}{\psi_{2D}} - \frac{1}{2} m (\omega_x^2 x^2 + \omega_y^2 y^2) - g \phi_0^2 |\psi_{2D}|^2 = -\frac{\hbar^2}{2m} \frac{\phi_0''}{\phi_0} + \frac{1}{2} m \omega_z^2 z^2$$

$$\Rightarrow \begin{cases} \frac{\hbar^2}{2m} \phi_0'' - \frac{1}{2} m \omega_z^2 z^2 \phi_0 + \frac{\hbar \gamma}{\epsilon} \phi_0 = 0 \Rightarrow -\frac{\hbar^2}{2m} \phi_0'' = \left[-\frac{1}{2} m \omega_z^2 z^2 + \gamma \right] \phi_0 \\ i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \psi + \left[\frac{1}{2} V_2 + g_3 \phi_0^2 / 4 \right] \psi \end{cases} \quad (2)$$

$$\boxed{\phi_0} \quad \phi_0(z) = A e^{-\frac{z^2}{2\beta^2}} \Rightarrow \cancel{\frac{\hbar^2}{2m} A \left(\frac{1}{4\beta^4} \right)} \phi_0' = A \left(-\frac{z}{\beta^2} \right) e^{-\frac{z^2}{2\beta^2}}$$

$$\Rightarrow \phi_0'' = A \left[-\frac{1}{\beta^2} e - \frac{z}{\beta^2} \left(-\frac{z}{\beta^2} \right) e \right]$$

$$\therefore \frac{\hbar^2}{2m} \left[-\frac{1}{\beta^2} + \frac{z^2}{\beta^4} \right] - \frac{m \omega_z^2 z^2}{2} + \frac{\hbar \gamma}{\epsilon} = 0$$

$$\Rightarrow \left[-\frac{\hbar^2}{2m\beta^2} + \frac{\hbar \gamma}{\epsilon} \right] + \left[\frac{\hbar^2}{2m\beta^4} - \frac{m \omega_z^2}{2} \right] z^2 = 0$$

$$\Rightarrow \gamma = \frac{\hbar^2}{2m\beta^2}, \quad m \omega_z^2 = \frac{\hbar^2}{m\beta^4}, \quad \gamma = \frac{\hbar^2}{2m\beta^2} \cdot \frac{1}{\beta^2} = \frac{\hbar^2}{2m a_z^2} = \frac{\hbar^2}{2m} \frac{m \omega_z^2}{\hbar} \quad \boxed{\gamma = \frac{\hbar \omega_z}{2}}$$

$$\Rightarrow \boxed{\beta = \frac{\hbar}{m \omega_z} = a_z} \quad \left[a_z \equiv \frac{\hbar}{m \omega_z} \right]$$

$$\therefore \phi_0(z) = A e^{-\frac{z^2}{2a_z^2}}$$

$$\int_{-\infty}^{\infty} e^{-\frac{z^2}{2a_z^2}} dz = \sqrt{2\pi} a_z$$

$$\therefore \text{if } \underbrace{\int |\phi_0|^2 = 1}_{\text{normalize}} \Rightarrow A^2 \int e^{-\frac{z^2}{a_z^2}} = 1 \Rightarrow A^2 a_z \sqrt{\pi} = 1$$

$$\Rightarrow \boxed{A = \pi^{-1/4} a_z^{-1/2}}$$

$$\therefore \boxed{\phi_0(z) = \pi^{-1/4} a_z^{-1/2} e^{-\frac{z^2}{2a_z^2}}} \Rightarrow \phi_0'' = \pi^{-1/4} a_z^{-1/2} \left[-\frac{1}{a_z^2} + \frac{z^2}{a_z^4} \right] e$$

$$= \left[-\frac{1}{a_z^2} + \frac{z^2}{a_z^4} \right] \phi_0$$

$$\boxed{\psi_2} \quad i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \psi + \frac{1}{2} V_2 +$$

$$GPE : i\hbar \psi_t = \left[-\frac{\hbar^2}{2m} \nabla^2 + \frac{1}{2} V_3 + g_3 |\psi|^2 \right] \psi$$

$$\psi_3 = \psi_2 \phi_0 e^{-i\frac{\delta t}{\hbar}}$$

Average over z

$$\int GPE \cdot \phi_0^* dz \Rightarrow i\hbar \frac{\partial \psi_2}{\partial t} \phi_0 + \delta \psi_2 \phi_0 = -\frac{\hbar^2}{2m} \nabla^2 \psi_2 \phi_0 - \frac{\hbar^2}{2m} \frac{\phi_0''}{\phi_0} \psi_2 + [V + g_3 \phi_0^2 |\psi_2|^2] \psi_2 \phi_0$$

$$\Rightarrow i\hbar \frac{\partial \psi_2}{\partial t} + \delta \psi_2 = -\frac{\hbar^2}{2m} \nabla^2 \psi_2 - \frac{\hbar^2}{2m} \left[-\frac{1}{a_z^2} + \frac{z^2}{a_z^4} \right] \psi_2 + [V + g_3]$$

$$\langle \phi_0^* (= \langle \phi_0) \Rightarrow i\hbar \frac{\partial \psi_2}{\partial t} \phi_0^2 + \delta \psi_2 \phi_0^2 = -\frac{\hbar^2}{2m} \nabla^2 \psi_2 \phi_0^2 - \frac{\hbar^2}{2m} \phi_0'' \phi_0 \psi_2 [\quad] \psi_2 \phi_0^2$$

$$\int dz \Rightarrow i\hbar \frac{\partial \psi_2}{\partial t} \phi_0^2 + \delta \psi_2 \phi_0^2 = -\frac{\hbar^2}{2m} \nabla^2 \psi_2 \phi_0^2 - \frac{\hbar^2}{2m} \left[-\frac{1}{2} m \omega_z^2 z^2 + \delta \right] \phi_0^2 \psi_2 [\quad] \psi_2 \phi_0^2$$

$$\int dz \Rightarrow i\hbar \frac{\partial \psi_2}{\partial t} + \delta \psi_2 = -\frac{\hbar^2}{2m} \nabla^2 \psi_2 - \frac{1}{2} m \omega_z^2 z^2 \psi_2 + \delta \psi_2 + [V_1 + V_2] \psi_2 + g_3 \phi_0^4 \psi_2$$

$[\phi_0^2 = 1]$

$$\Rightarrow i\hbar \frac{\partial \psi_2}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \psi_2 + V_2 \psi_2 + g_3 \int \phi_0^4 |\psi_2|^2 \psi_2$$

$$\therefore g_2 \equiv g_3 \int \phi_0^4 = g_3 \frac{1}{2} \frac{\sqrt{2}}{\pi} \frac{1}{a_z} = g_3 \frac{1}{\sqrt{2\pi} a_z}$$

$$\therefore \boxed{g_2 = g_3 \frac{1}{\sqrt{2\pi} a_z}}$$

$$\Rightarrow \boxed{i\hbar \frac{\partial \psi_2}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \psi_2 + \left[\frac{1}{2} m (\omega_x^2 x^2 + \omega_y^2 y^2) + g_2 |\psi_2|^2 \right] \psi_2} \quad (2)$$

$$\psi_3 = \psi_2 \phi_0 e^{-i\frac{\delta t}{\hbar}}, \quad \phi_0(z) = \pi^{-1/4} a_z^{-1/2} e^{-\frac{z^2}{2a_z^2}},$$

$$a_z^2 \equiv \frac{\hbar}{m\omega_z}, \quad \delta \equiv \frac{\hbar\omega_z}{2} \Rightarrow \frac{\delta}{\hbar} = \frac{\omega_z}{2} \Rightarrow [\quad] = \frac{1}{2} \checkmark$$

$$g_2 \equiv g_3 \frac{1}{\sqrt{2\pi} a_z}, \quad g_3 = \frac{4\pi \hbar^2 a^2}{m} \Rightarrow \sqrt{2\pi} a_z g_2 = \frac{4\pi \hbar^2 a^2}{m}$$

Adm:

$$\omega_r \equiv \omega_x = \omega_y$$

$$\tau = \omega_z t, \quad R = r/\lambda, \quad \zeta^2 = \frac{\hbar}{m\omega_z}$$

$$\begin{matrix} (2) \\ \xrightarrow{t \rightarrow \tau, R \rightarrow R} \end{matrix}$$

$$i\hbar\omega_z \frac{\partial \Psi}{\partial \tau} = -\frac{\hbar^2}{2m} \frac{m\omega_z}{\hbar} \nabla_R^2 \Psi + \left[\frac{1}{2} m\omega_r^2 \frac{\hbar}{m\omega_z} R^2 + g_2 |\Psi|^2 \right] \Psi$$

$$\div \hbar\omega_z \quad i\Psi_\tau = -\frac{1}{2} \nabla_R^2 \Psi + \left[\frac{1}{2} \Omega^2 R^2 + \frac{g_2}{\hbar\omega_z} |\Psi|^2 \right] \Psi$$

$$\boxed{\Omega \equiv \frac{\omega_r}{\omega_z}}$$

$$\bar{\Psi} = (4\pi\zeta^2 a)^{1/2} \Psi = \left(\frac{\hbar}{m\omega_z} \frac{g_3 m}{4\pi\hbar} \right)^{1/2} \Psi = \left(\frac{g_3}{\omega_z} \right)^{1/2} \Psi$$

$$\Rightarrow i\bar{\Psi}_\tau = -\frac{1}{2} \nabla_R^2 \bar{\Psi} + \left[\frac{1}{2} \Omega^2 R^2 + \frac{g_3}{\sqrt{4\pi} a} \frac{1}{\hbar\omega_z} \frac{\omega_z}{g_3} |\bar{\Psi}|^2 \right] \bar{\Psi}$$

$$\text{Adm } \Psi: \quad b\bar{\Psi} = \Psi \Rightarrow i\bar{\Psi}_\tau = -\frac{1}{2} \nabla^2 \bar{\Psi} + \left[\frac{1}{2} \Omega^2 R^2 + \frac{g_2}{\hbar\omega_z} b^2 |\bar{\Psi}|^2 \right] \bar{\Psi}$$

$$\therefore b^2 = \frac{\hbar\omega_z}{g_2}$$

$$\Rightarrow \boxed{\Psi_2 = \sqrt{\frac{\hbar\omega_z}{g_2}} \bar{\Psi}_2}$$

$$\left[\begin{aligned} b^2 &= \frac{\hbar\omega_z}{g_3} = \frac{\hbar\omega_z m}{4\pi\hbar^2 a} = \frac{m\omega_z}{4\pi a \hbar} \\ 4\pi\zeta^2 a &= 4\pi \frac{\hbar}{m\omega_z} a \end{aligned} \right]$$

$$\begin{aligned} &= \boxed{\frac{\partial \bar{\Psi}_2}{\partial \tau} = -\frac{1}{2} \nabla^2 \bar{\Psi}_2 + \left[\frac{1}{2} \Omega^2 R^2 + |\bar{\Psi}_2|^2 \right] \bar{\Psi}_2} \\ &\quad \Psi_3 = \Psi_2 \phi, \quad \phi(z) = \pi^{-1/4} a_z^{-1/2} e^{-\frac{z^2}{2a_z^2}} \\ &\quad \zeta^2 = a_z^2 \equiv \frac{\hbar}{m\omega_z}, \quad \delta \equiv \frac{\hbar\omega_z}{2}, \quad \Omega = \omega_r/\omega_z \\ &\quad \Psi_2 = \sqrt{\frac{\hbar\omega_z}{g_2}} \bar{\Psi}_2, \quad g_2 \equiv g_3 \frac{1}{\sqrt{4\pi} a_z}, \quad g_3 = \frac{4\pi\hbar^2 a}{m} \\ &\quad \tau = \omega_z t, \quad R = r/\lambda, \quad \zeta^2 = \frac{\hbar}{m\omega_z} = a_z^2 \end{aligned}$$

Adimensionalized

+

Reduced 3D-2D

GPE.